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# SIZE EFFECTS ON MAGNETIC PROPERTY AND CRYSTAL STRUCTURE OF $Mn_3O_4$ NANOPARTICLES IN MESOPOROUS SILICA

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$Mn_3O_4$  nanoparticles with particle sizes of 7.8, 11.4, and 18.3 nm were synthesized in the pores of mesoporous silica, and their crystal structure and magnetic properties were investigated. The powder X-ray diffractions at room temperature indicated that the crystal structural symmetry was the same as that for bulk crystal, and the lattice constants deviated from those for bulk crystal, which depended on the particle size. In addition, compared with the bulk crystal, the Jahn-Teller distortion for the nanoparticles was suppressed and decreased with decreasing the particle size. The coercive field for 7.8 nm was rather smaller than those for 11.4 and 18.3 nm. The nanoparticles with 11.4 and 18.3 nm exhibited pronounced three kinds of magnetic transition temperatures, whereas the susceptibility for 7.8 nm indicated the existence of two transition temperatures. These experimental results suggested that the  $Mn_3O_4$  nanoparticles have a strong correlation between crystallographic structure and magnetic property, and the characteristic magnetic size effects are attributed to the reduction of Jahn-Teller distortion.

*Crystal structure, Lattice distortion, magnetic property, Nanoparticles*

## I. INTRODUCTION

Spinel oxides with a general formula  $AB_2O_4$  has tetrahedral A site with four oxygen and octahedral B site with six oxygen arranged in a pyrochlore lattice. When magnetic ions occupy the B site, geometric magnetic frustration often occurs owing to the strong antiferromagnetic interactions between B sites. The  $Mn_3O_4$  has geometrical frustration derived from the antiferromagnetic interactions between the B sites of spinel oxide.  $Mn_3O_4$  exhibits interesting phenomena such as magnetodielectric, magnetoelastic and magnetocaloric behaviors, and successive phase transformations owing to the strong coupling among spin, orbital, and lattice degrees of freedom [1-6]. For  $Mn_3O_4$ , the tetrahedral A sites and octahedral B sites are selectively occupied by  $Mn^{2+}$  and  $Mn^{3+}$  ions, respectively. Since  $Mn^{3+}$  ions at the  $MnO_6$  octahedral site with one electron in the doubly degenerate  $e_g$  states,  $Mn_3O_4$  exhibits a structural phase transition from cubic to tetragonal at 1443 K due to the strong instability of the Jahn-Teller distortion, whose magnitude is estimated to be  $c/\sqrt{2}a \sim 1.16$  based on the ratio of lattice constants ( $a \sim 5.76$  Å and  $c \sim 9.47$  Å) [4,5]. The  $Mn_3O_4$  bulk crystal exhibits a ferrimagnetic ordering, called Yafet-Kittel phase, at  $T_N = 42$  K, where the magnetic moments are oriented along the [110] direction. In addition, as temperature decreases, the magnetic phase of  $Mn_3O_4$  changes successively to an incommensurate (IC) phase at  $T_1 = 40$  K, and finally to a commensurate phase, called cell-doubled (CD) phase at  $T_2 = 34$  K, where the magnetic unit cell is doubled along the [110] direction from that of the original cubic spinel structure.

The nanoparticle of strongly correlated materials such as manganese oxide are expected to exhibit characteristic size effects on crystal structure and magnetic property owing to the strongly electron correlation and strong coupling among spin, orbital, and lattice. In fact, we have already shown that nanoparticles of strongly correlation materials,  $LaMnO_3$  [7],  $DyMnO_3$  [8], and  $NiO$  [9], exhibit characteristic size dependences on the magnetic properties and crystallographic

structures. These nanoparticles exhibit the remarkable changes in crystallographic structure, i.e., changes in the structural symmetry, lattice constant, and lattice strain, which induces the unique size effects of magnetic properties. Moreover, the nanoparticles of the magnetically frustrated materials are expected to the characteristic size effects. The nanoparticles of  $DyMn_2O_5$ , which has magnetic frustration owing to the intricate interplay of at least five magnetic Mn-Mn spin interactions in the geometrically frustrated structure, exhibit the unusual magnetic size effect derived from the distortion on the crystallographic structure [10]. The magnetic properties for the  $Mn_3O_4$  nanoparticles are expected to be different from those for bulk crystal owing to the modulation of magnetic frustration and lattice distortion due to the surface effect. In this study, we synthesized the  $Mn_3O_4$  nanoparticles with particle size of approximately 8-18 nm in the pores of mesoporous silica, and their crystallographic structure and magnetic properties were investigated. These results suggested the existence of the strong correlation between magnetic properties and lattice distortion associated with Jahn-Teller distortion in the  $Mn_3O_4$  nanoparticles.

## II. EXPERIMENT

$Mn_3O_4$  nanoparticles were synthesized in the pores, with a diameter of approximately 7 nm, of mesoporous silica SBA-15. SBA-15 was used as a template to equalize the particle size in the fabrication of the  $Mn_3O_4$  nanoparticles and to deflocculate synthesized nanoparticles since it has a two-dimensional hexagonal mesoporous structure and its pores are separated by silica wall [11]. The  $Mn_3O_4$  nanoparticles were synthesized by soaking the SBA-15 in a aqueous solution of  $MnCl_2 \cdot 4H_2O$ . The soaked SBA-15 was then dried and calcinated in an air atmosphere. The powder X-ray diffraction (XRD) measurements for the synthesized nanoparticles were carried out at room temperature by using a synchrotron radiation X-ray diffractometer with a Debye-Scherrer camera at the beamline BL-8B at the Photon Factory (PF) of the Institute of

Materials Structure Science, High Energy Accelerator Research Organization (KEK). The X-ray wavelength was calibrated using the XRD pattern of CeO<sub>2</sub> powder. The magnetic properties of the nanoparticles were investigated by using a superconducting quantum interference device (SQUID) magnetometer (Quantum Design Inc.). Temperature dependence of DC susceptibility and magnetization curves were measured.

### III. EXPERIMENTAL RESULTS AND DISCUSSION

We observed powder XRD patterns for the synthesized Mn<sub>3</sub>O<sub>4</sub> nanoparticles in the pores of SBA-15 at room temperature, as shown in Fig. 1. The diffraction patterns of three Mn<sub>3</sub>O<sub>4</sub> nanoparticles exhibited broad Bragg peaks, which were attributed to the tetragonal symmetry with space group *I4<sub>1</sub>/amd* the same as that of bulk crystal. Note that the particle sizes of the Mn<sub>3</sub>O<sub>4</sub> nanoparticles were estimated based on the peak positions and the full widths at half maximum of the Bragg peaks using Scherrer's equation. These results indicated successful synthesis the Mn<sub>3</sub>O<sub>4</sub> nanoparticles with mean particle size of  $d = 7.8, 11.4,$  and  $18.3$  nm.

**FIG. 1 HERE** (Note white space above and below.)

The lattice constants  $a$  and  $c$  of the nanoparticles were estimated from XRD patterns. The observed broad diffraction peaks were deconvoluted to those associated with Mn<sub>3</sub>O<sub>4</sub> by multi-peak fitting. The lattice constants were calculated from the relation between the lattice constants and the plane indices determined from the Bragg peak angles for all the deconvoluted peaks. The estimated lattice constants are listed in Table I. The lattice constants for the nanoparticles were slightly different from those for bulk crystal [4] and depended on the particle size. Compared with the lattice constants for bulk crystal, the deviation of value of  $c$  is larger than that of  $a$ , which indicated that the crystallographic structure for the nanoparticle is distorted from the bulk crystal. In particular, for  $d = 7.8$  and  $11.4$  nm, the lattice constant  $a$  is larger than that for bulk crystal, while  $c$  is smaller than that for bulk crystal. The lattice constant  $a$  for  $d = 7.8$  nm was slightly larger than that for  $d = 11.4$  and  $18.3$  nm, while the  $c$  for  $d = 7.8$  and  $11.4$  nm were similar value and slightly larger than that for  $d = 18.3$  nm. The results indicate that the lattice distortion for the Mn<sub>3</sub>O<sub>4</sub> nanoparticles increases as particle size decreases. The tetragonal distortion has been used to describe the Jahn-Teller distortion in Mn<sub>3</sub>O<sub>4</sub>. The Mn<sub>3</sub>O<sub>4</sub> bulk crystal has huge tetragonal distortion,  $c/(\sqrt{2} a) = 1.162$  [4]. The estimated tetragonal distortion for the nanoparticles listed in Table I indicates that the Jahn-Teller distortion for Mn<sub>3</sub>O<sub>4</sub> nanoparticles is suppressed as compared with the bulk crystal, and its magnitude decreases with decreasing particle size.

**TABLE 1 HERE** (Note white space above and below.)

The DC magnetic susceptibilities for the nanoparticles were measured under an external magnetic field  $H = 200$  Oe in both field cooling (FC) and zero-field-cooling (ZFC) conditions. Fig. 2(a) shows the temperature dependence of the Mn<sub>3</sub>O<sub>4</sub> nanoparticles with  $d = 11.4$  nm. The susceptibility exhibited a clear hysteresis between FC and ZFC conditions. Both the FC and ZFC susceptibilities exhibited a rapid increased at around  $T = 42$  K, corresponding to the  $T_N$  for bulk crystal, which is the onset of the Yafet-Kittel phase. In addition, the ZFC susceptibility exhibited the two anomalies at around 40 K as shown in inset of Fig. 2 (a), which correspond to the transition temperatures  $T_1$  and  $T_2$  as described later. Fig. 2(b) shows temperature dependence of FC susceptibilities for the three nanoparticles, which suggest that the  $T_N$  of 7.8 nm is lower than those of 11.4 and 18.3nm. The magnetic ordering temperatures were estimated by temperature differentiation of ZFC susceptibility as shown in Fig. 2(c) [6,12]. The  $d\chi_{ZFC}/dT$  curves for  $d = 11.4$  nm shows three peaks corresponding to the  $T_N = 42.4$  K,  $T_1 = 40.9$  K, and  $T_2 = 39.3$  K, respectively. The evaluated magnetic transition temperatures are listed in Table II. Compared with the bulk crystal [6], the  $T_N$  and  $T_1$  are almost the same as those of bulk crystal, while the  $T_2$  is slightly higher than that of bulk crystal. For  $d = 18.3$  nm, the  $d\chi_{ZFC}/dT$  curves also shows three peaks, and those transition temperatures are almost same values as those for  $d = 11.4$  nm. However, the  $d\chi_{ZFC}/dT$  curves for  $d = 7.8$  nm exhibit distinct two peaks at 40 K corresponding to the  $T_N$  and 34.5 K corresponding to the  $T_1$  or  $T_2$ , which suggests the disappearance of the IC or CD phase for nanoparticles with  $d = 7.8$  nm. In the case of Mn<sub>3</sub>O<sub>4</sub> thin films [12], the  $dM/dT$  curves of thickness of 8 and 25 nm exhibit three peaks corresponding to the three transition temperatures, while that of thickness of 4 nm exhibits a peak only. The transition temperature for 4 nm was larger than that for 8 and 25 nm. The authors concluded that the behavior for thin film with thickness of 4 nm was associated with the surface roughness, which could greatly destroy the exchange interactions among the Mn ions. Hence, the  $T_N$  for the thinnest is lower than that of the thicker film [12]. In this study, it is considered that the decrease in the transition temperatures for the nanoparticles with  $d = 7.8$  nm is attributed to the surface state.

**FIG. 2 HERE** (Note white space above and below.)

The magnetization curves for the nanoparticles were measured ranging from 5 to -5 T. Fig. 3 shows the magnetization curves for the nanoparticles at 5 K. The magnetization curves exhibited a hysteresis loop below  $T_N$ . The values of coercive field  $H_c$  at 5 K are listed in Table II. The  $H_c$  for the  $d = 11.4$  and  $18.3$  nm were almost same value, which was rather larger than that for 7.8 nm. The obtained  $H_c$  is quite larger than that for powdered bulk polycrystalline,  $H_c = 4980$  Oe. This behavior is similar to previous studies for the nanoparticle [13,14] and thin film [12,15]. In the case of the previous investigations for Mn<sub>3</sub>O<sub>4</sub> nanoparticles, it is

concluded that the increases in the coercive field are attributed to the significant increase in specific surface, such that the magnetic surface anisotropy is enhanced [13,14].

**FIG. 3 HERE** (Note white space above and below.)

The crystal structural analysis results indicated that the lattice constants and Jahn-Teller distortion depended on the particle size and exhibited the pronounced change at 7.8 nm, while those values were almost constant above 11.4 nm. The nanoparticles for  $d = 11.4$  and 18.3 nm exhibited the three pronounced magnetic transition temperatures, while the nanoparticles for  $d = 7.8$  nm exhibited two transition temperatures. In addition, the values of the transition temperatures and the coercive field for  $d = 7.8$  nm was rather lower and smaller than those for  $d = 11.4$  and 18.3 nm, respectively. These results indicated that the magnetic properties for the  $\text{Mn}_3\text{O}_4$  nanoparticles exhibited pronounced changes below 7.8 nm. The experimental results suggest the  $\text{Mn}_3\text{O}_4$  nanoparticles have strong correlation between crystal structure and magnetic property. As particle size decreases, the changes in the crystal structure and electronic state should be induced because of the enhancement of influence of surface state, which results in the remarkable changes in the lattice constants, Jahn-Teller distortion, and magnetic properties for the nanoparticles of  $d = 7.8$  nm. The transition temperature  $T_2$  for the nanoparticles is higher than that for bulk crystal. It is considered that the increases in the transition temperatures are attributed to the decreasing of atomic distances and enhancing the exchange interactions of Mn atoms owing to the reduction of the lattice constants [12]. On the other hand, the  $\text{Mn}_3\text{O}_4$  has a tetragonal distortion combined with Jahn-Teller distortion. The  $\text{Mn}^{3+}$  spins at the  $\text{MnO}_6$  octahedra having antiferromagnetic interactions with each other on the corner of a tetragonally distorted tetrahedron give rise to two possible energetically degenerate spin configurations, resulting in the increase in the spin fluctuation above  $T_N$  and its decrease with successive magnetic phase transitions [3]. In the case of the nanoparticles of  $\text{DyMnO}_3$  with orthorhombic distortion of perovskite structure associated with the Jahn-Teller distortion, the size dependences of magnetic properties are related to the reduction of the lattice distortion [8]. Moreover, the  $\text{Mn}_3\text{O}_4$  bulk crystal shows magnetic field induced crystal structural- and magnetostructural-modulation and magnetic phase transition [1,2]. The structural phase transition is associated with a rearrangement of frustrated one-dimensional antiferromagnetic chains composed by B-site  $\text{Mn}^{3+}$  [2]. In addition, the structural- and magnetic- phase transition is also induced under high pressure for bulk crystal and nanoparticle [1,5,16]. Compared with bulk crystal, the  $\text{Mn}_3\text{O}_4$  nanoparticles with  $d \sim 10$  nm show an elevation of phase transition pressure and different phase diagram, which are derived from unique atomic configuration and higher surface energy together with the nano-size effect [5]. In this study, the decrease in the Jahn-Teller distortion for  $d = 7.8$  nm is presumed to induce the

change in the orbital degeneracy of the  $e_g$  states in  $\text{Mn}^{3+}$  ions. Therefore, the modulation of the lattice constants and reduction of the Jahn-Teller distortion varies the magnetic interactions between the Mn spins, which result in the pronounced changes in the transition temperatures, coercive field and magnetic phase diagram. By continuing statistical research on  $\text{Mn}_3\text{O}_4$  nanoparticles with other particle sizes, we are planning to understand the size effect of  $\text{Mn}_3\text{O}_4$  nanoparticles in more detail.

**TABLE 2 HERE** (Note white space above and below.)

#### IV. CONCLUSION

We synthesized the  $\text{Mn}_3\text{O}_4$  nanoparticles with  $d = 7.8$ , 11.4, and 18.3 nm in the pores of mesoporous silica SBA-15 and investigated their crystallographic structure and magnetic property. The XRD patterns of the three nanoparticles exhibited the tetragonal symmetry similarly to the bulk crystal at room temperature. The lattice constants for the nanoparticles were slightly different from those for bulk crystal and depended on particle size. The lattice constants change significantly at 7.8 nm, in addition, the tetragonal distortion associated with Jahn-Teller distortion decreases at 7.8 nm. The magnetic properties exhibited pronounced changes at 7.8 nm. The temperature dependences of DC magnetic susceptibility exhibited the successive phase transformations. The nanoparticles with  $d = 11.4$  and 18.3 nm clearly exhibited the three transition temperatures, whereas the nanoparticles with  $d = 7.8$  nm exhibited the two transition temperatures. The magnetization curve for the nanoparticles exhibited a hysteresis loop at below  $T_N$ . The coercive field for  $d = 11.4$  and 18.3 nm was almost the same value, and its value for  $d = 7.8$  nm decreased significantly. These experimental results suggested that the  $\text{Mn}_3\text{O}_4$  nanoparticles with  $d = 7.8$ -18.3 nm have a strong correlation between the crystallographic structure and magnetic properties. The reduction of the tetragonal lattice distortion associated with Jahn-Teller distortion results in the changes in the magnetic transition temperature, coercive field, and magnetic phase diagram.

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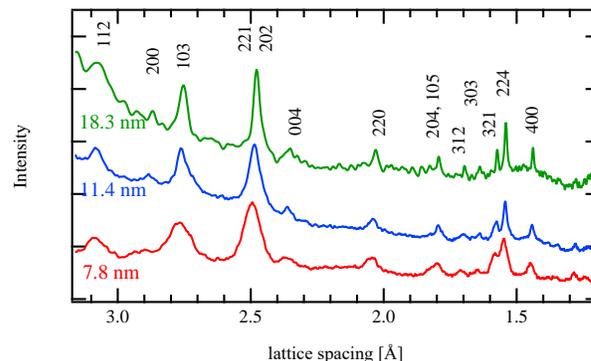


Fig. 1. XRD patterns for the  $Mn_3O_4$  nanoparticles with  $d = 7.8, 11.4,$  and  $18.3$  nm at room temperature.

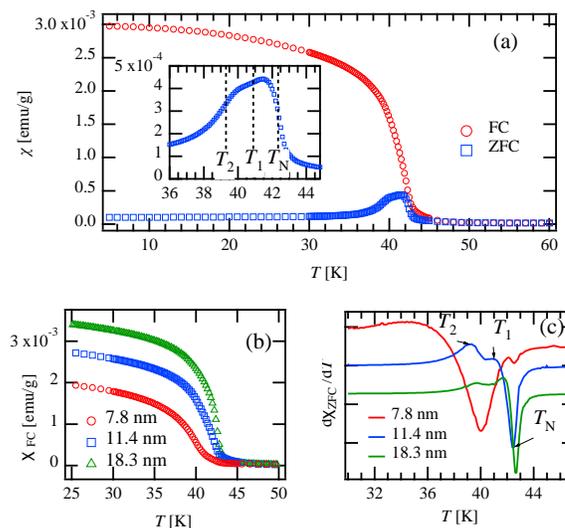


Fig. 2. (a) Temperature dependence of DC susceptibility for the nanoparticles with  $d = 11.4$  nm under magnetic field  $H = 200$  Oe in FC and ZFC conditions. Inset shows a magnification of ZFC susceptibility near the transition temperatures. (b) Temperature dependence of susceptibility for  $d = 7.8, 11.4,$  and  $18.3$  nm under FC condition. (c) Temperature dependence of derivative curve  $d\chi_{ZFC}/dT$ .

TABLE I

LATTICE CONSTANTS AND TETRAGONAL DISTORTION

$d$ [nm]	$a$ [Å]	$c$ [Å]	$c/\sqrt{2}a$
7.8	5.780(1)	9.448(5)	1.156
11.4	5.765(1)	9.445(1)	1.159
18.3	5.762(1)	9.432(1)	1.158
bulk [4]	5.76289(1)	9.46885(1)	1.162

Lattice constants,  $a$  and  $c$ , of tetragonal symmetry and tetragonal distortion,  $c/\sqrt{2}a$ , for the nanoparticles and bulk crystal at room temperature. The values for bulk crystal are referred from Ref. 4.

TABLE II

MAGNETIC TRANSITION TEMPERATURES AND COERCIVE FIELD

$d$ [nm]	$T_N$ [K]	$T_1$ [K]	$T_2$ [K]	$H_c$ [Oe]
7.8	40.0		34.5	8190
11.4	42.4	40.9	39.3	8740
18.3	42.6	41.7	39.7	8790
bulk	43.9 [6]	41.2 [6]	35.0 [6]	4980

Magnetic transition temperatures and coercive field at 5 K for the nanoparticles and bulk crystal. The transition temperatures is estimated by temperature differentiation of ZFC susceptibility. For  $d = 7.8$  nm, the lower transition temperature is corresponding to the  $T_1$  or  $T_2$ . The transition temperatures for the bulk crystal are referred from Ref. 6.

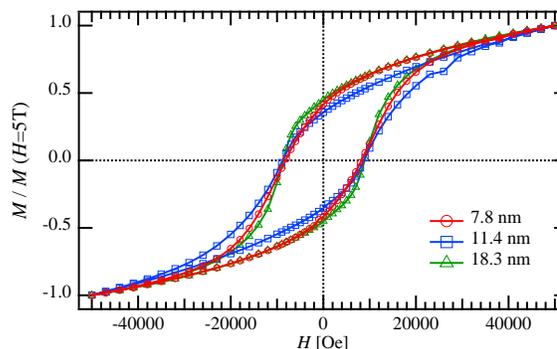


Fig. 3. Magnetization curves for the  $Mn_3O_4$  nanoparticles at 5 K.