

Intermittent dynamics of antiferromagnetic phase in inhomogeneous iron-based chalcogenide superconductor

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The coexistence of phases, characterized by different electronic degrees of freedom, commonly occurs in layered superconductors. Among them, alkaline intercalated chalcogenides are model systems showing the microscale coexistence of paramagnetic (PAR) and antiferromagnetic (AFM) phases, however, the temporal behavior of different phases is still unknown. Here, we report a visualization of the atomic motion in the granular phase of $K_xFe_{2-y}Se_2$ using x-ray photon correlation spectroscopy. Unlike the PAR phase, the AFM texture reveals an intermittent dynamics with avalanches as in martensites. When cooled down across the superconducting transition temperature T_c , the AFM phase goes through an anomalous slowing behavior, suggesting a direct relationship between the atomic motions in the AFM phase and the superconductivity. In addition to providing compelling evidence of avalanchelike dynamics in a layered superconductor, the results provide a basis for different theoretical models to describe the quantum states in inhomogeneous solids.

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I. INTRODUCTION

The observation of superconductivity in iron-based chalcogenides [1] has opened new frontiers in the field of layered materials with an interesting interplay between atomic defects, magnetism, and superconductivity [2]. Such an interplay has been widely discussed for other layered systems [3–7]. Among iron-based chalcogenides, the $A_xFe_{2-y}Se_2$ ($A = K, Rb, Cs$) system [8–11] is a good example, showing an intrinsic phase separation [12–18] and a delicate balance between a magnetic phase due to an iron-vacancy order and the coexisting metallic phase. $A_xFe_{2-y}Se_2$ shows superconductivity below a transition temperature T_c of ~ 32 K and manifests a peculiar microstructure with a coexisting antiferromagnetic (AFM) phase having the stoichiometry of $A_{0.8}Fe_{1.6}Se_2$ (245) and a paramagnetic (PAR) metallic phase of $A_xFe_2Se_2$ (122). A variety of experiments have studied the phase separation properties [8–18], revealing a wealth of information on the microstructure of the system. For example, space-resolved micro-x-ray diffraction (μ XRD) on $K_xFe_{2-y}Se_2$ [12] has identified a $\sqrt{5} \times \sqrt{5}$ superstructure due to iron-vacancy order in the average tetragonal lattice to occur below ~ 580 K and a phase separation to appear below ~ 520 K. The former is a second-order transition while the latter transition has primarily a first-order character [19]. Depending on the growth conditions, the system contains about

$\sim 70\%$ – 90% of an insulating AFM phase with a $\sqrt{5} \times \sqrt{5}$ superstructure while the remaining minority phase is metallic and is characterized as a compressed in-plane lattice. This peculiar phase separation puts these chalcogenides in the class of granular systems in which the dynamics in the microscopic granules has a large effect on their macroscopic properties. Here, we have used x-ray photon correlation spectroscopy (XPCS), a diffraction-based technique [20–22], to probe the atomic dynamics in the coexisting phases of superconducting $K_xFe_{2-y}Se_2$. XPCS exploits the temporal evolution of the x-ray speckle pattern generated by coherent radiation. The speckle patterns represent a direct fingerprint of the nanoscale phase disorder in the material. If the material fluctuates in time, the speckle pattern does the same, and a measurement of the speckle intensity fluctuation reveals the dynamics of the system.

II. EXPERIMENTAL DETAILS

Single-crystal samples of $K_xFe_{2-y}Se_2$ were prepared using the Bridgman method [9]. After the growth, the single crystals were sealed into a quartz tube and annealed for 12 h at 600°C . A well-characterized sample of size $3 \times 3 \text{ mm}^2$, having the composition $K_{0.65}Fe_{1.65}Se_2$, was used for the present measurements. The electric and magnetic characterizations were

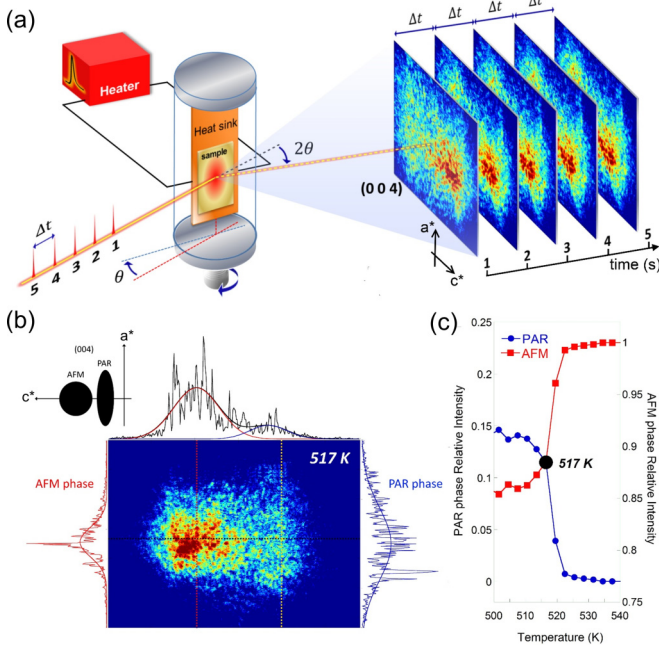


FIG. 1. X-ray photon correlation spectroscopy (XPCS) measurements on $K_xFe_{2-y}Se_2$. (a) Schematic diagram of the experimental setup with the sample mounted on a copper block inside an evacuated heat chamber. (b) A typical speckle pattern of the (004) reflection at 517 K, showing two different phases characterized by their c axis, i.e., expanded and compressed c axis for AFM and PAR phases, respectively (see the cartoon picture). Line profiles of the intensity distributions of the speckles corresponding to AFM (left) and PAR (right) phases along with the averaged profiles are also shown. The upper profile shows the two phases characterized by a different c axis. (c) Temperature evolution of the AFM (red squares) and PAR (blue dots) phases across the phase separation temperature ~ 520 K.

performed by temperature-dependent measurements of resistivity using a physical property measurement system (PPMS, Quantum Design) and magnetization using a superconducting quantum interference device (SQUID) magnetometer (Quantum Design). The sample exhibits a sharp superconducting transition at T_c of ~ 32 K.

The XPCS experiments were carried out in the $\theta/2\theta$ reflection geometry with a beam falling parallel to the b axis of the single-crystal sample having tetragonal symmetry [see, e.g., Fig. 1(a) showing the experimental geometry]. The measurements were carried out at the Coherence Beamline P10 of PETRA III synchrotron radiation source in Hamburg where the x-ray beam, produced by a 5-m-long undulator (U29), is monochromatized using a Si(111) double-crystal monochromator. An x-ray photon beam of 8 keV energy with a bandwidth $dE/E \sim 1.4 \times 10^{-4}$ was used. At this energy the transverse coherence length is $277 \mu\text{m}$ in the vertical direction and $46 \mu\text{m}$ in the horizontal direction. The collimated coherent x-ray beam was focused using a beryllium compound refractive lens (CRL) translocator [23] to a size of about $2 \times 2 \mu\text{m}^2$ on the sample positioned 1.6 m downstream of the translocator center. The incident flux on the sample was $\sim 10^{11}$ photons/s. The exit window of the heating chamber and He cryostat (see, e.g., the Supplemental Material [24]

and Ref. [25] therein) as well as the entrance window of the detector flight path was covered by a $25\text{-}\mu\text{m}$ -thick Kapton sheet. The scattered signal was detected at a distance of ~ 5 m using a large horizontal scattering setup. A PILATUS 300 detector (7 ms readout time) was used for the alignment and a MAXIPIX 2×2 detector (0.3 ms readout time) was used to record the x rays scattered by the sample with an angular resolution of 6.228×10^{-4} deg.

III. RESULTS AND DISCUSSIONS

Figure 1(a) shows a schematic picture of the XPCS setup. The single-crystal sample of $K_xFe_{2-y}Se_2$ is mounted on a copper block inside an evacuated chamber. More details on the sample environment and the experimental setup are shown in the Supplemental Material [24] (see also Ref. [25] therein). The sample shows a phase separation while it is cooled across a temperature of ~ 520 K. A speckled (004) Bragg reflection, measured on the $K_xFe_{2-y}Se_2$ crystal at a constant temperature of 517 K, is displayed in Fig. 1(b). The reflection is a direct indicator of the phase separation [12,19] in the block antiferromagnetic tetragonal phase due to the iron-vacancy order (space group $I4/mmm$ with $a = b = 4.01 \text{ \AA}$, $c = 13.84 \text{ \AA}$, hereafter called the AFM phase) and the c -axis expanded tetragonal paramagnetic phase (hereafter called the PAR phase). The profiles of (004) reflections are shown in Fig. 1(b) displaying the typical speckles due to coherent x rays. The temperature dependence of the normalized intensity for the two phases is shown in Fig. 1(c). The majority AFM phase contributes $\sim 80\%$ – 90% while the remaining $\sim 10\%$ – 20% is the PAR phase.

After 100 s of measurements at 517 K, the sample temperature was raised quickly by 1 K to bring the system into a nonequilibrium state. Figure 2 displays the time evolution of the two phases before and after a temperature step of 1 K. The relaxation can be seen in Fig. 2(a) displaying the time evolution of the integrated intensities corresponding to the two phases (normalized with respect to the total intensity in the equilibrium state). The time evolution of the (004) reflection mean profiles for the two phases is shown in Fig. 2 of the Supplemental Material [24]. The sample temperature was kept constant (at 518 ± 0.2 K) during the whole time series measurements [see, e.g., the temporal fluctuations of temperature plotted in Fig. 2(a)]. The speckled pattern evolves with time after the temperature stimulation (i.e., a quick change of temperature by 1 K after ~ 100 s), shown for different instants in Fig. 2(b). There are some apparent changes as a function of time in the two regions of the speckled pattern while the system is relaxing from the nonequilibrium state.

The complex nonequilibrium dynamics in highly heterogeneous systems can be visualized in the best way through two-time correlation functions (ttcf) [26–31]. The two-time correlation functions are calculated by correlating all possible pairs of diffraction patterns collected during the time series described above. The following equation is used to calculate the ttcf = $C(I(t_1), I(t_2))$ [32],

$$\frac{\sum_m [I_m(t_1) - \langle I(t_1) \rangle][I_m(t_2) - \langle I(t_2) \rangle]}{\sqrt{\sum_m [I_m(t_1) - \langle I(t_1) \rangle]^2 [I_m(t_2) - \langle I(t_2) \rangle]^2}}. \quad (1)$$

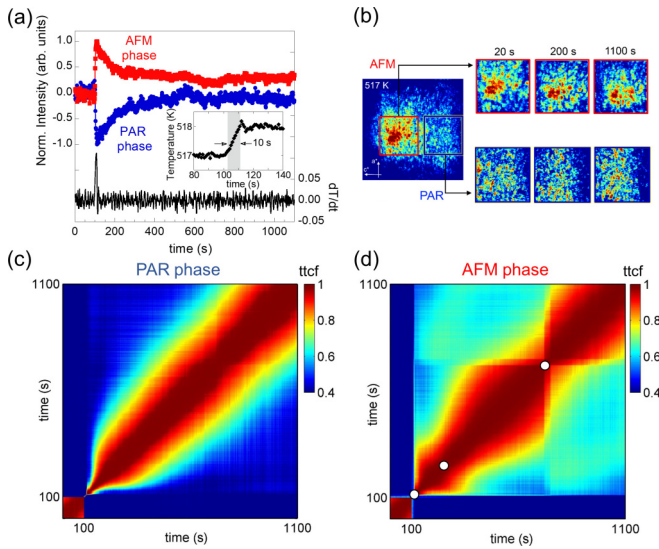


FIG. 2. Time evolution of different phases before and after the temperature pulsed step of 1 K from the equilibrium state at 517 K. (a) Time evolution of the AFM (red squares) and PAR (blue dots) phases. The time series collection started at a constant temperature (517 K) in the equilibrium state. After 100 s the temperature was changed rapidly by 1 K (black curve shown as the inset) and 1000 additional diffraction patterns were collected to study the complex nonequilibrium dynamic. (b) Speckle patterns for some time delays during the time series for the AFM (upper panels) and PAR (lower panels) phases. (c), (d) Two-time correlation functions. (c) The PAR phase shows a normal (quasistatic) dynamics while the situation for (d) the AFM phase is different, revealing avalanches in the domain transformations.

Here, $C(I(t_1), I(t_2))$ is the two-time correlation function (ttcf), $I_m(t_1)$ and $I_m(t_2)$ are intensities measured in the detector pixel m at time t_1 and t_2 , and $\langle I(t_1) \rangle$ and $\langle I(t_2) \rangle$ are respectively the mean intensities measured over all pixels of images recorded at time t_1 and t_2 . To minimize the overlap, we tried different regions of interest around (004) reflections corresponding to the two phases in $K_x\text{Fe}_{2-y}\text{Se}_2$ before judging for the regions of interest shown as squares in Fig. 2(b) for the calculations of ttcf. The two-time correlation images for the two phases are displayed in Figs. 2(c) and 2(d). The PAR phase displays a normal (quasistatic) dynamics revealing the system to evolve from a locked-static state to the next one defined by a close minimum in the energy landscape [33–35]. In fact, the intensity distribution in the two-time correlation image is spreading out with time. The diagonal width of the two-time correlation image provides information on the correlation time, i.e., the timescale in which the given atomic configuration (characterized by a well-defined wave vector) no longer corresponds to the one measured at a later time. The width of the intensity distribution in the ttcf image for the PAR phase increases with time, indicating increased correlation times with time, i.e., the slowing down of the dynamics of this phase with time (for a system in a dynamic equilibrium, the width of the two-time correlation image is expected to be a constant, i.e., no time evolution). This indicates the formation of larger and larger domains of the PAR phase at the expense of small domains, consistent with models for growth processes [31].

In the potential energy landscape approach [33], the dynamics of the PAR phase suggests that the system finds itself in a configuration space with a deep energy basin and evolving towards a deeper and deeper local energy minima.

Unlike the PAR phase, the AFM phase shows a very peculiar dynamics. Indeed, the two-time correlation of the AFM phase [Fig. 2(d)] reveals dramatic decorrelation events characterized by the sudden narrowing of the intensity distribution profile appearing intermittently with time. This temporal intermittence indicates an avalanchelike atomic dynamics in the majority AFM phase. In fact, such an intermittent dynamics describes rearrangements to localized microcollapses of groups of particles, which trigger subsequent collapses in the neighboring regions through the formation of stress dipoles. Therefore, the intermittent progression of the AFM phase can be identified as an incubation time effect, i.e., silent growth and explosions in sequences. This avalanchelike dynamics has been found in a number of physical phenomena including martensitic transformations [36,37], deformation of metallic glasses [35], crystallization of a hard-sphere glass [38], and shear flow of droplet emulsions through a thin opening [39]. It is likely that these events of microscopic rearrangement act as important mediators in the particular phase via the cooperative relief of atomic-level strain between the coexisting AFM and PAR phases. This particular phase is the interface (INT) phase identified in a microdiffraction study on the same material [40]. Therefore, the AFM and PAR phases are separated by a well-defined INT phase. Thus, the intermittent dynamics of the AFM texture in $K_x\text{Fe}_{2-y}\text{Se}_2$ is intrinsic and indicative of a complex energy landscape with numerous minima (different equilibrium states) in which the system stays for long periods of time in stable configurations, reflecting both localized and cascade relaxation dynamics.

After the study of the nonequilibrium dynamics in which the sample temperature was varied sharply by 1 K, the sample was kept at a constant temperature (at 518 K) for a long time (more than 1 h). Assuming the sample to be in the equilibrium state, we measured a second time series, collecting speckle patterns for 500 s. The instantaneous autocorrelation function $g_2(t)$ [20–22] was calculated using this time series revealing characteristic correlation times (τ) to be ~ 550 and ~ 400 s, respectively, for the PAR and AFM phases (see Supplemental Material for a detailed description [24] and Ref. [41] therein).

The fact that the AFM phase in $K_x\text{Fe}_{2-y}\text{Se}_2$ shows avalanchelike and intermittent temporal fluctuations in the collective dynamics poses a question if such a dynamics has any relationship with the superconductivity in $K_x\text{Fe}_{2-y}\text{Se}_2$. To search for a possible connection between the dynamics and the superconductivity, we have studied the speckle evolution for the two phases while the sample is cooled down across the superconducting transition temperature T_c . For this purpose, we have varied the temperature with a constant rate. The temperature evolution of the mean diffraction profiles is shown in Fig. 4 of the Supplemental Material [24]. Here, the dynamics has been studied by evaluating the temperature-temperature correlation function (TTcf). The TTcf has been calculated using the same equation used for the calculations of the ttcf [Eq. (1)] in which the time is replaced by temperature. The procedure has been commonly used to explore the response of atomic dynamics across temperature-dependent transitions [32]. Figures 3(a) and 3(b) show the TTcf for

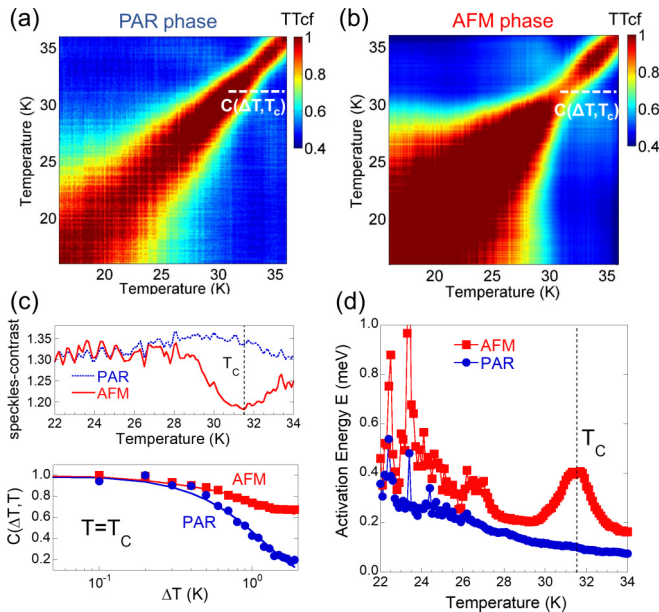


FIG. 3. Nanodomain dynamics across superconducting transition temperature (T_c). (a), (b) Two-temperature correlation function (TTcf) images calculated for a time series collected during a linear temperature ramp. The TTcf of the PAR phase (a) shows a normal cone shape indicating a linear acceleration of domain dynamics with increasing temperature. Instead, the TTcf of the AFM phase (b) shows a clear anomaly around T_c . (c) Temperature evolution of the speckle contrast (upper panel) for PAR (dotted blue line) and AFM (solid red line) phases. The lower panel shows the autocorrelation function $C(\Delta T, T)$ at $T = T_c$ for PAR (blue dots) and AFM (red squares) phases. The $C(\Delta T, T)$ curves have been extracted selecting horizontal cuts from the TTcf images of the two phases and normalized to the speckle contrast. $C(\Delta T, T)$ curves are fitted using the stretched exponential behavior [Eq. (2)]. (d) The activation energy (E) extracted from the fits is plotted as a function of temperature for the PAR (blue dots) and AFM (red squares) phases. The activation energy shows a bump at $T = T_c$ for the AFM phase.

the PAR and AFM phases. Apparently, both phases display similar dynamics upon cooling, however, a clear anomaly for the AFM phase around T_c can be seen. At this temperature the intensity distribution is sharper before the spreading out, i.e., the correlation time at T_c is much smaller for the AFM phase. This indicates large fluctuations near T_c followed by a slowing of the AFM phase below the transition temperature. On the other hand, the PAR phase seems to evolve normally. Therefore, the AFM phase shows anomalous dynamic correlations across T_c in the phase-separated $K_x\text{Fe}_{2-y}\text{Se}_2$.

To have detailed insight a further analysis of the temperature-dependent speckle patterns was done. The autocorrelation function $C(\Delta T, T)$ was calculated at different temperatures in the shown interval around T_c . The following equation was used to describe the calculated autocorrelation functions,

$$C(\Delta T, T) = A \exp \left[- \left(\frac{\Delta T}{E} \right)^\beta \right], \quad (2)$$

where A represents the speckle contrast, β is the shape parameter of the stretched exponential function, and E is the

activation energy. The evaluated speckle contrast around T_c for the two phases is displayed in Fig. 3(c) (upper). In the lower panel of Fig. 3(c) we have shown the $C(\Delta T, T)$ at T_c and is normalized with respect to the speckle contrast. The activation energy (E) around T_c is also shown in Fig. 3(d). It is evident from Fig. 3(c) (top) that the speckle contrast around T_c is much smaller for the AFM phase than for the PAR phase, implying the existence of faster fluctuations in the AFM phase at T_c . The speckle contrast at T_c does not drop to zero, meaning that the dynamics is not fully decorrelating. On the other hand, the autocorrelation $C(\Delta T, T)$ at T_c reveals that the AFM phase has a slower relaxation than the PAR phase [see Fig. 3(c), bottom] but with an anomalously increased activation energy [Fig. 3(d)]. These observations indicate that there should be some other much slower processes actively incorporated within the AFM phase. Therefore, the dynamics of the AFM phase is indeed more complex with at least two relaxation channels present which are well separated in activation energy. The underlying correlation function could resemble that of glasses or other disordered systems with coexisting fast and slow relaxations.

IV. CONCLUSIONS

In summary, we have studied the dynamics of nanodomains in a phase-separated $K_x\text{Fe}_{2-y}\text{Se}_2$ system. While the minority PAR phase reveals the commonly known steady slowing down with time, the majority AFM phase shows intermittent nonequilibrium dynamics as in martensites involving cooperative atomic rearrangements with avalanches. This complex dynamics of the AFM phase may have some direct correlation with the superconductivity in $K_x\text{Fe}_{2-y}\text{Se}_2$. Indeed, the measurements across the superconducting transition temperature show that the AFM phase goes through an anomalous atomic dynamics across the superconducting transition temperature, reflecting the involvement of complex energy landscapes to establish the superconducting quantum state. It is worth recalling that superconductivity is accompanied by a hardening in local atomic modes, that has been seen in a series of superconducting families [42–44]. Therefore, the behavior of the AFM phase could be a result of the superconducting transition that can be mediated by lattice fluctuations (Fe-Fe lattice) or spin fluctuations in the AFM phase. The fact that the local magnetic moment decreases sharply at T_c [shown by x-ray emission spectroscopy (XES) [45]] as well the AFM order tends to suppress at T_c [46] makes it plausible to think that the superconductivity in these materials may have some exotic mechanism involving the collective mode of lattice and spin characterized by slow dynamics. It should be mentioned that, in addition to the PAR and AFM phases, the system is characterized by the INT phase which forms out of the AFM phase [40,45]. Therefore, it is likely that the superconductivity appears in the INT phase as argued earlier. However, more efforts are required to clarify issues on the dynamics of the INT phase.

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