

# Soft x-ray irradiation induced metallization of layered TiNCl

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## Abstract

We have performed soft x-ray spectroscopy in order to study the photoirradiation time dependence of the valence band structure and chemical states of layered transition metal nitride chloride TiNCl. Under the soft x-ray irradiation, the intensities of the states near the Fermi level ( $E_F$ ) and the  $Ti^{3+}$  component increased, while the Cl  $2p$  intensity decreased. Ti  $2p$ - $3d$  resonance photoemission spectroscopy confirmed a distinctive Fermi edge with Ti  $3d$  character. These results indicate the photo-induced metallization originates from deintercalation due to Cl desorption, and thus provide a new carrier doping method that controls the conducting properties of TiNCl.

## Introduction

TiNCl is an interesting layered material that belongs to a family of layered nitride halides  $MNX$  ( $M = Ti, Zr, Hf$  and  $X = Cl, Br, I$ ) systems with two polymorphs, the FeOCl-type crystal structure ( $\alpha$ -form) and the SmSI-type crystal structure ( $\beta$ -form) [1, 2]. TiNCl is known to take only the FeOCl-type crystal structure [2], consisting of a stack of orthogonal  $M$ - $N$  layers located between two  $X$  layers. TiNCl is a band insulator with a direct band gap of approximately 0.5-3 eV [3, 4], and theoretical studies have proposed its application to optoelectronic devices [4], photocatalysts [5], and spin devices [6]. Upon Na intercalation between TiNCl layers, TiNCl becomes a superconductor with a relatively high superconducting transition temperature of  $\sim 16$  K [7]. As electron-doped

1  $\beta$ -form HfNCl and ZrNCl [8-20], electron-doped TiNCl is considered a candidate for  
2 unconventional superconductors, where exotic mediation forces for Cooper pairing other  
3 than phonon are discussed [21-25].

4 In order to further explore the physical properties of this remarkable material, it is  
5 essential to control carrier concentrations. In the  $\beta$ -form  $MNX$ , the carrier control has  
6 been performed by intercalation of an alkali- or alkaline-earth-metal [8, 9, 15, 16], and/or  
7 off stoichiometry or deintercalation of  $X$  [26]. As for TiNCl, to the best of our knowledge,  
8 carrier control by neither element substitution nor off stoichiometry has been reported.  
9 This may be due to the fact that TiNCl is easily thermally decomposed into TiN by  
10 annealing[27-29]. The intercalation of alkali metal atoms and/or organic molecules  
11 between layers is the only method way for doping carriers. However, intercalated samples  
12 are unstable in a humid air, and precise carrier controlling is challenging, which prevents  
13 the systematic investigations of physical properties systematically.

14 One of the techniques to create conductive samples is photo-irradiation-induce  
15 metallization by breaking chemical bonds [30-36]. In SrTiO<sub>3</sub> and TiO<sub>2</sub>, a metallic state  
16 called two-dimensional electron gases is created on the surface by photoirradiation. Two-  
17 dimensional electron gases have attracted extensive attention owing to unique physical  
18 properties [30-35, 37-39], providing opportunities for the development of next generation  
19 of electronic and photonic devices. If photo-induced bond breaking is also effective in  
20 TiNCl, it can be established as a new method to control the physical properties of TiNCl.  
21 Therefore, it is one of the important challenges to investigate whether TiNCl is metallized  
22 by photoirradiation.

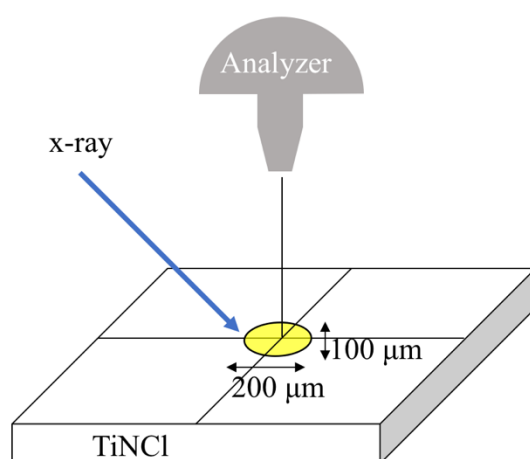
23 In this paper, we performed valence and core level photoelectron spectroscopy (PES),  
24 x-ray absorption spectroscopy (XAS), and Ti  $2p$ - $3d$  resonance photoelectron  
25 spectroscopy (RPES) of layered transition metal nitride chloride TiNCl. It was found that  
26 soft x-ray irradiation induced the metallization of TiNCl, and the metallization was  
27 closely related to the Cl desorption. This is the first report of electron-doped TiNCl due  
28 to Cl desorption.

## 30 **Experiment**

31 The highly crystalline TiNCl was grown by a method described elsewhere [7]. It was  
32 pelletized into disk shape in an Ar filled glove box. Subsequently, the pellet was  
33 transferred to an ultra-high vacuum chamber for photoelectron spectroscopy without  
34 being exposed to the humid air for a long time. The experiments were performed at the  
35 soft x-ray beam line BL23SU of SPring-8 [40] using a photoemission spectrometer  
36 equipped with a Gammadata-Scienta SES-2002 electron analyzer. In the RPES

1 measurement, the energy resolution was set to 80 meV, and in the other PES  
2 measurements, it was set to less than 200 meV. To reduce a possibility for pyrolysis of  
3 TiNCl, we cooled down the sample and measured it at 100 K instead of at 300 K. We  
4 performed PES measurements using a photon energy of 1000 eV. For calibration of the  
5 binding energy, we used the Fermi edge of a gold film, which was located close to the  
6 sample and had a good electrical contact to the sample.

7 The irradiation light of  $h\nu = 1000$  eV was used, and the detector angle was set to  $90^\circ$ .  
8 Sets of Ti  $2p$ , N  $1s$ , Cl  $2p$ , and valence spectra were also measured at  $h\nu = 1000$  eV in the  
9 time interval of 25 minutes without changing the measured location of the sample surface.  
10 The irradiation time includes the time of the PES measurements. Though it took 8 min to  
11 take a set of Ti, N, Cl, and valence spectra, we refer to the irradiation time as the time  
12 when we started the measurement of Ti. The spot size of the excitation light was  $100 \times$   
13  $200 \mu\text{m}^2$ , and the photon flux was  $2 \times 10^{12}$  Photons/sec [40]. The illustration of the  
14 experimental geometry is shown in Fig. 1. The intensity of PES spectra was normalized  
15 by the area intensity of Ti  $2p$  core-level spectra, while those of RPES were normalized by  
16 the photocurrent. XAS spectrum was measured using the total electron yield (TEY) mode.  
17 To obtain a clean surface, the TiNCl pellet was fractured at 100 K under an ultrahigh  
18 vacuum of  $1.5 \times 10^{-8}$  Pa.



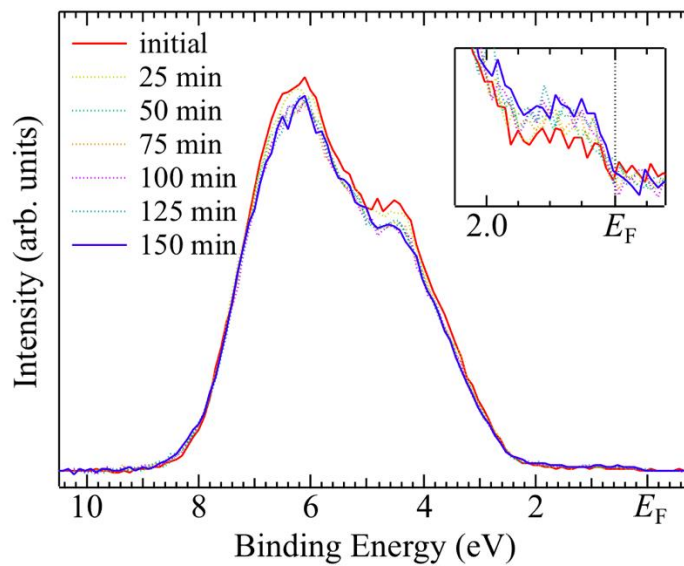
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20 Figure 1. Schematic diagram of the x-ray irradiation.  
21

## 22 Results and discussion

23 Figure 2 shows the soft x-ray irradiation time dependence of the valence band PES

1 spectra of TiNCl. The background of an iterative Shirley method [41, 42] was subtracted  
 2 from the raw spectra. The initial valence band spectrum of TiNCl had a peak at 6 eV with  
 3 a shoulder structure at 4 eV, and there was almost no intensity in the near- $E_F$  region. The  
 4 overall spectral shape is consistent with that of the previous study [25], and thus the states  
 5 of approximately 6 and 4 eV can be ascribed to Cl 3*p* and N 2*p* orbitals, respectively.

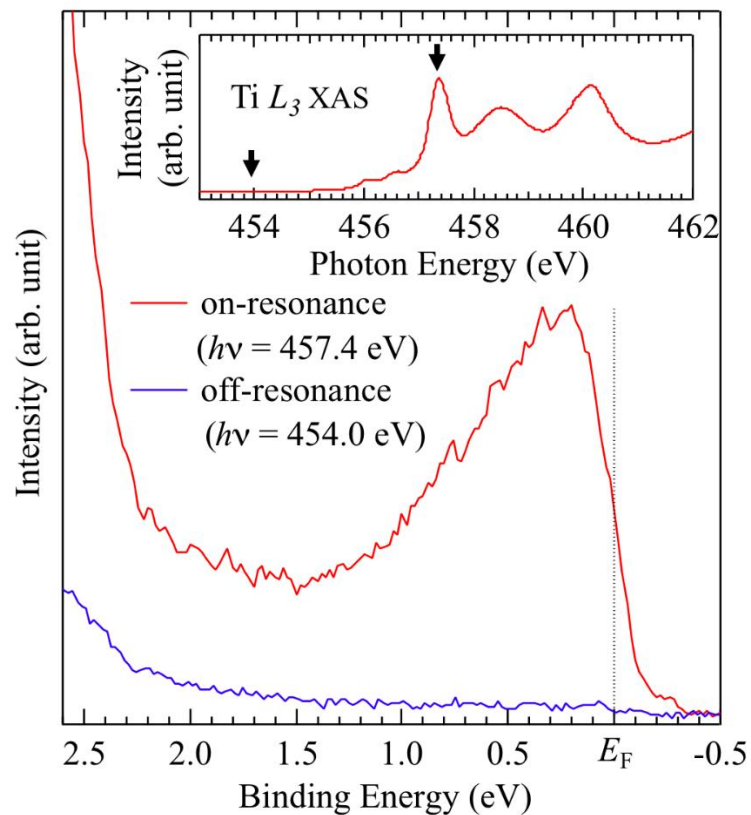
6 The peak intensities at 6 eV and 4 eV were decreased gradually. In contrast, the  
 7 spectral intensity in the region between  $E_F$  and  $\sim 1.5$  eV appears to be increased, as  
 8 illustrated in the inset that shows a blowup of the spectra near  $E_F$ . The intensity near  $E_F$   
 9 of the initial spectrum can be ascribed to the states induced by the soft x-ray irradiation  
 10 during the measurement and/or may be assigned to impurity states. In the spectrum after  
 11 150 min irradiation, a clear Fermi edge structure was observed, suggesting the metallic  
 12 nature of the measured region of the sample surface. The change in the valence band and  
 13 the intensity near  $E_F$  due to the soft x-ray irradiation is very similar to the change due to  
 14 the photoirradiation effect of SrTiO<sub>3</sub> [30].



15  
 16 Figure 2. Soft x-ray irradiation time dependence of the valence band PES spectra of TiNCl,  
 17 measured with a photon energy of 1000 eV at 100 K. Inset shows the enlargement of the  
 18 spectra near  $E_F$ .

19  
 20 To spectroscopically confirm the metallization and the character of the states at  $E_F$ ,  
 21 we performed RPES of irradiated TiNCl at the Ti 2*p*-3*d* absorption threshold. The photon  
 22 energies for the RPES measurement were chosen based on the Ti  $L_3$  XAS spectrum, as  
 23 shown in the inset of Fig. 3. Ti  $L_3$  XAS spectrum had sharp absorption peaks at 457.4 eV,  
 24 458.5 eV, and 460.1 eV in the region of Ti 2*p*<sub>3/2</sub>-3*d* absorption edge. In the octahedral

1 structure, the  $3d$  level was split into  $e_g$  and  $t_{2g}$  states by the crystal field effect. However,  
 2 since the Ti ions in TiNCl are in a distorted octahedral structure [7], the degeneracy was  
 3 further removed by the static Jahn-Teller effect, leading to the observed three-peak  
 4 structure. We measured RPES spectra of TiNCl in the off-resonance ( $h\nu = 454.0$  eV) and  
 5 on-resonance ( $h\nu = 457.4$  eV) conditions after the irradiation the soft x-rays for 4 h. While  
 6 the spectral intensity at  $E_F$  in the off-resonance spectrum was almost negligible, the one  
 7 in the on-resonance spectrum exhibited a distinctive Fermi edge, and was considerably  
 8 enhanced. The resonance enhancement of the near- $E_F$  structure indicates that its orbital  
 9 character is Ti  $3d$ . Furthermore, the presence of the distinctive Fermi edge indicates that  
 10 the soft x-ray irradiation indeed induces the metallization of TiNCl.

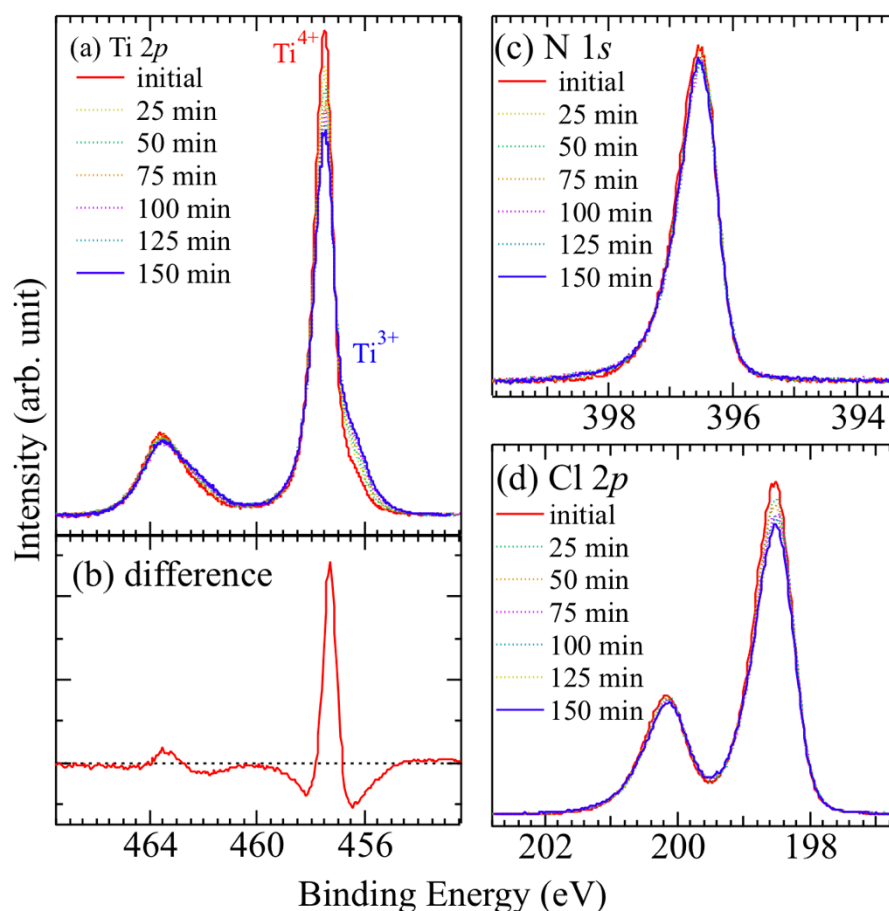


11  
 12 Figure 3. ON (red) and OFF (blue) Ti  $2p$ - $3d$  resonance photoemission spectra of the  
 13 irradiated TiNCl measured with a photon energy of 457.4 eV and 454.0 eV, respectively.  
 14 The inset shows the Ti  $2p$  absorption spectrum where the employed photon energies for  
 15 the resonance photoemission are indicated by arrows.

16  
 17 To understand the mechanism of the metallization induced by soft x-ray irradiation in  
 18 terms of chemical states, we have measured the core level spectra of TiNCl with different  
 19 irradiation conditions. Figure 4 (a) shows soft x-ray irradiation time dependence of the Ti  
 20  $2p$  core-level spectra, and Fig.4 (b) represents the initial and 150 min difference spectrum.

1 The peaks at 457.5 eV and 463.5 eV are spin-orbit split Ti 2*p*, namely Ti 2*p*<sub>3/2</sub> and Ti 2*p*<sub>1/2</sub>,  
2 respectively [43]. The initial Ti 2*p* spectrum showed a sharp Ti<sup>4+</sup> peak at 457.5 eV [25]  
3 and a small shoulder structure at 456.5 eV. In the Ti 2*p* spectrum, as the irradiation time  
4 increased, the peak area of Ti<sup>4+</sup> decreased compared to the initial value. In addition, the  
5 intensity of the shoulder structure at 456.5 eV clearly increased gradually. This binding  
6 energy is different from that of TiN (455.2 ± 0.2 eV) [44, 45], but is very similar to the  
7 weak Ti 2*p* spectral component of electron-doped SrTiO<sub>3</sub> and electron-doped anatase-  
8 TiO<sub>2</sub>, which is assigned to the component of Ti<sup>3+</sup> state [30, 31, 46-48]. Thus, the  
9 observation indicates that the soft x-ray irradiation on TiNCl induces the changes Ti<sup>4+</sup> in  
10 the valence state of the Ti ion from the Ti<sup>3+</sup> state to the Ti<sup>4+</sup> state. We performed spectral  
11 fitting using two components. The area intensities of Ti<sup>3+</sup> and Ti<sup>4+</sup> components at each  
12 irradiation time are summarized in Table 1.

13 Figure 4 (c) compares the spectra of the N 1*s* core level at each soft x-ray irradiation  
14 time. These spectra had peaks at 396.5 eV, and were almost identical, suggesting that  
15 there was no significant change in the spectral shape due to soft x-ray irradiation. The Cl  
16 2*p* core-level spectra consist of Cl 2*p*<sub>3/2</sub> (198.4 eV) and Cl 2*p*<sub>1/2</sub> (200.1 eV) spin-orbit split  
17 peaks. The intensity of the Cl 2*p* spectrum after the soft x-ray irradiation of 150 min  
18 decreased gradually, which suggests the desorption of Cl atoms from the surface is  
19 induced by the soft x-ray irradiation. This is consistent with the observed reduction of the  
20 valence band spectra, which reflects a dominant Cl 3*p* contribution.



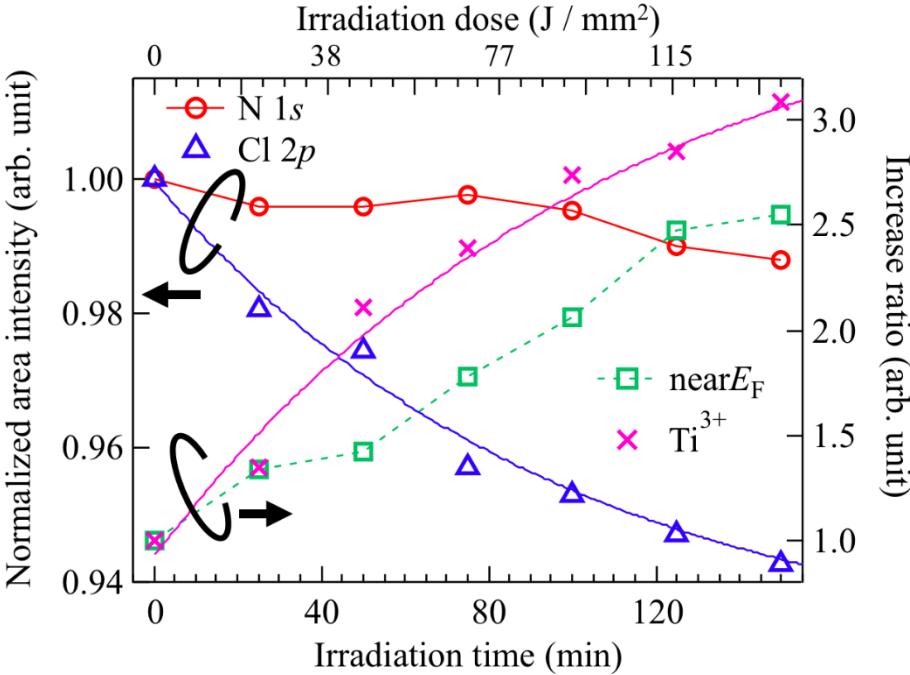
1  
 2 Figure 4. Soft x-ray irradiation time dependence of core-levels PES spectra. (a) Ti  $2p$   
 3 core-level spectra irradiated from the initial to 150 min in the interval of 25 min. (b)  
 4 Difference between the initial and 150min. (c) and (d) are analogous spectra to (a) for the  
 5 N  $1s$  and Cl  $2p$  core-level, respectively.  
 6  
 7 Table 1. Area intensity ratio estimated from  $Ti^{3+}$  and  $Ti^{4+}$  in  $Ti\ 2p$  core level PES spectra  
 8 at each irradiation time.

Irradiation time (min)	Relative area intensity of components (%)	
	$Ti^{4+}$ in $Ti\ 2p$	$Ti^{3+}$ in $Ti\ 2p$
0	90.60	9.40
25	87.37	12.63
50	80.22	19.78
75	77.52	22.48
100	74.26	25.74
125	73.21	26.79
150	70.96	29.04

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Here, we discuss the relationship between the observed results. Figure 5 shows the soft x-ray irradiation time dependence of the intensity of N 1s and Cl 2p core-level PES spectra, the PES intensity between  $E_F$  and 1.5 eV, and the  $Ti^{3+}$  ratio ( $Ti^{3+} / (Ti^{4+} + Ti^{3+})$ ) of Ti 2p core-level PES spectrum. We normalized their initial values to unity. The intensity of N 1s core-level PES does not show a marked change depending on the soft x-ray irradiation time of soft x-ray. In a comparison with that of the N 1s core-level PES, the intensity of the Cl 2p core-level PES decreased more rapidly with increasing irradiation time. Furthermore, simultaneously with a decrease in Cl, the areas of the  $Ti^{3+}$  component and of the states near  $E_F$  increased in a similar manner. A smaller increase in near  $E_F$  intensity may be attributed to a more complicated orbital character of states near  $E_F$  than that of the  $Ti^{3+}$  core level. These observations suggest the close relationship of the desorption of Cl atoms and an increase in the intensity near  $E_F$  with an increase in the  $Ti^{3+}$  component. We consider that the desorption of a portion of the Cl atoms from the TiNCl layers introduces electron carriers into the layers, making the layer conductive, which is inferred from the observation of the metallic Fermi edge. Therefore, to estimate the time of change, the data for  $Ti^{3+}$  and Cl 2p, which are more intense than the near- $E_F$  data, were fitted with an exponential function ( $A\exp(-t/\lambda)+B$  for  $Ti^{3+}$  and  $C\exp(t/\lambda)+D$  for Cl 2p, where  $t$  and  $\lambda$  are time and lifetime, respectively.). The obtained lifetimes of  $Ti^{3+}$  and Cl 2p were  $119 \pm 29$  and  $96 \pm 31$  min, respectively. This provided a value of 100 min as the order estimate of the time of change for this experiment.



22



1 Figure 5. Soft x-ray irradiation time dependence of the core-level and near- $E_F$  intensities  
2 of TiNCl. Area intensity of N 1s and Cl 2p core-level spectra (left axis), and an increase  
3 ratio in area intensity of near  $E_F$  and  $Ti^{3+}$  in Ti 2p core level spectra (right axis). The initial  
4 values were normalized to 1.

5  
6 There are two possible causes of the desorption of Cl atoms. One of them is a  
7 pyrolysis. TiNCl has been reported to be decomposed completely into TiN when it is  
8 heated above 550°C [27-29]. In the present experiment, however, the sample was kept at  
9 100 K during the measurement, and therefore this possibility can be ruled out. The other  
10 possible cause is photon stimulated desorption [49-51]. There are primarily two  
11 mechanisms for photon stimulated desorption, which are distinguished by the process  
12 induced by photon irradiation: The Menzel, Gomer, and Redhead (MGR) model and the  
13 Knotek and Feibelman (KF) model [49, 50]. The MGR model involves the excitation of  
14 valence electrons and works reasonably well for covalent bonds. The KF model involves  
15 the excitation of core electrons and is applicable to ion bonding. Since the TiNCl is  
16 considered to be an ionic insulator composed of  $TiN^+$  and  $Cl^-$  [23], we assumed that this  
17 mechanism is involved in the Cl desorption of TiNCl. In the case of  $TiO_2$ , photoinduced  
18 desorption is explained by the KF model in terms of Coulomb expulsion between  $Ti^{4+}$   
19 and  $O^+$  that are induced by emitting Auger electrons after the relaxation of electrons in O  
20 2p levels to create photo-holes in the Ti 3p levels [48]. If the KF mechanism is involved  
21 in the present case, the electron transfer that leads to Coulomb expulsion may occur from  
22 Cl 3p (valence band) to Ti 3p. Further experiments are needed to confirm the KF  
23 mechanism, as it is we think beyond the scope of the present study.

24 With photon stimulated Cl desorption, the final products may be TiN or Cl deficient  
25  $TiNCl_{1-\delta}$ , depending on the amount of Cl deficiency. However, from the Ti 2p  
26 photoelectron spectrum, no sharp TiN peak was observed at  $455\pm 0.3$  eV[44, 45],  
27 suggesting the formation of  $TiNCl_{1-\delta}$ , rather than TiN. In  $\beta$ -(Hf or Zr) NCl, Cl deficiency  
28 is called deintercalation, which is one of the methods of electron doping to the system.  
29 Deintercalated  $HfNCl_{0.7}$  is a metallic system that also exhibits superconductivity below  
30 24 K [26]. However, since electron doping utilizing the Cl deintercalation into  $\alpha$ - $MNX$   
31 has never been reported, the present study is the first experimental realization of  
32 deintercalated electron doped  $\alpha$ - $MNX$ . Further Cl deintercalation by the soft x-ray  
33 irradiation would make the system superconductive, like  $\beta$ -form compounds.

34 Lastly, we discuss the benefits of methodologies and findings. Regarding pyrolysis,  
35 bond breaking by heat cannot be confined within a small area because of the diffusion of  
36 heat in a solid. In addition, temperature controlling of a sample is severe for a target

1 material located in a narrow temperature region of a complicated phase diagram. In  
2 contrast, bond breaking by light irradiation occurs only in the place where light is  
3 irradiated. In particular, when core-levels excitation are important in the bond-breaking  
4 process as in the KF model, element/chemical site-selective bond breaking is possible by  
5 setting the excitation energy to a specific element/chemical site [52, 53]. This site-  
6 selective bond breaking is called a molecular scalpel and is mainly studied in surface  
7 adsorption molecular systems. Such a characteristic, combined with sub-micrometer scale  
8 spot sizes in third-generation synchrotron facilities, might be used to make sub-micron  
9 scale conducting paths at any place on the surface. Photo-induced metallization of TiNCl  
10 has enabled detailed electronic-structure investigation studies of metallic (even  
11 superconductive) TiNCl by controlling the carrier concentration in a small single crystal  
12 region on the surface, which shed light on the exotic properties of TiNCl.

## 14 **Conclusion**

15 In summary, the irradiation time dependence of the valence band and core-level  
16 spectra of TiNCl was studied by PES and Ti  $2p$ - $3d$  RPES. Using soft x-ray irradiation,  
17 the intensity of the valence band structure origination from Cl  $3p$  and N  $2p$  states  
18 decreased while the intensity near  $E_F$  increased. The intensity near  $E_F$  exhibited resonant  
19 enhancement, and the Fermi edge was clearly observed by Ti  $2p$ - $3d$  RPES. The analyses  
20 of core-level and near- $E_F$  PES spectral intensities as a function of the soft x-ray irradiation  
21 time revealed that the desorption of Cl atoms, the increase in the intensity of the  $Ti^{3+}$   
22 component, and increase of the spectral weight in the vicinity of  $E_F$  exhibit strong  
23 correlations. These results indicate that photoirradiation induced metallization of TiNCl  
24 occurs due to electron doping through Cl desorption. As estimated from Cl  $2p$  and  $Ti^{3+}$   
25 data, the order of the time-of-changes is 100 min. This technique can be used to further  
26 explore the conducting properties, especially the unconventional superconductivity, of  
27 TiNCl.

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36 BL25SU (Ref. 25). We assume that there are three reasons for this difference:

1 charging up, difference in metallization, and calibration of the spectrometer. Since  
2 the photon flux density of BL25SU is higher than that of BL23SU, larger binding  
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